# Abundance analysis of APOGEE spectra for 58 metal-poor stars from the bulge spheroid 

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#### Abstract

The central part of the Galaxy host a multitude of stellar populations, including the spheroidal bulge stars, stars moved to the bulge through secular evolution of the bar, inner halo, inner thick disk, inner thin disk, as well as debris from past accretion events. We identified a sample of 58 candidate stars belonging to the stellar population of the spheroidal bulge, and analyse their abundances. The present calculations of $\mathrm{Mg}, \mathrm{Ca}$, and Si lines are in agreement with the APOGEE-ASPCAP abundances, whereas abundances of $\mathrm{C}, \mathrm{N}, \mathrm{O}$, and Ce are re-examined. We find normal $\alpha$-element enhancements in oxygen, similar to magnesium, Si , and Ca abundances, which are typical of other bulge stars surveyed in the optical in Baade's Window. The enhancement of $[\mathrm{O} / \mathrm{Fe}]$ in these stars suggests that they do not belong to accreted debris. No spread in N abundances is found, and none of the sample stars is N -rich, indicating that these stars are not second generation stars originated in globular clusters. Ce instead is enhanced in the sample stars, which points to an s-process origin such as due to enrichment from early generations of massive fast rotating stars, the so-called spinstars.


Key words: Abundances - Atmospheres - Galaxy Bulge

## 1 INTRODUCTION

The stellar populations in the central part of the Galaxy can inform us about its complex formation processes. This region was recently

[^1]confirmed to contain stars in a metal-poor spheroidal bulge (e.g. Babusiaux et al. 2010, Dékány et al. 2013, Babusiaux 2016, Zoccali et al. 2018, Savino et al. 2020, Kunder et al. 2020, Arentsen et al. 2020, Queiroz et al. 2021 and references therein), along with a metal-rich contribution from the bar and inner thin disk, thick disk and halo interlopers. In addition, debris of past accretion events, such
as Gaia-Enceladus-Sausage (GES) (Belokurov et al. 2018, Helmi et al. 2018), and many other dwarf galaxy remnants, and minor substructures, absorbed during the early stages of the Galaxy formation (see e.g. Fernández-Trincado et al. 2022, Horta et al. 2020, 2021, 2022) are present. Therefore studies of the Galactic bulge region are important for understanding the early stages of our Galaxy's formation (e.g., Barbuy et al. 2018a, Rojas-Arriagada et al. 2020). In particular, Queiroz et al. $(2020,2021)$ combining distance derivation with proper motions from the Gaia Early Data Release 3 (Gaia Collaboration Brown et al. 2021) revealed stars of large eccentricity, but with orbits confined to the bulge region - with a maximum height from the Galactic mid-plane, $|\mathrm{z}|_{\text {max }}$, below 3 kpc , with intermediate metallicities, which are good candidates for belonging to the oldest Galactic bulge component (which we here call spheroid bulge stars).

The spheroidal metal-poor bulge can be thought of as a pressure supported structure formed through violent processes, such as hierarchical clustering via minor mergers, at a very early stage of the Galaxy. Ferraro et al. (2021) finds evidence that clumps of stars and gas existed at the time of the Milky Way formation. N-body simulations assume instead that early stellar discs heat rapidly as they form, and can lead to different density distributions for metal-rich and metal-poor stars (e.g. Debattista et al. 2017). Many other options are possible to form the metal-poor spheroid such as a major merger, accretion of dwarf galaxies, among others (e.g. Barbuy et al. 2018a). Whatever process, it leads to an observed metal-poor spheroid, and it has also to explain the very old ages of the in-situ globular clusters such as e.g. HP 1 (Kerber et al. 2019), Djorgovski 2 (Ortolani et al. 2019), Palomar 6 (Souza et al. 2021), of ages derived to be of $12.8 \pm 0.9,12.7 \pm 0.7$, and $12.4 \pm 0.9 \mathrm{Gyr}$, respectively.

The search for the earliest stars in the Galaxy is an important endeavour to try to identify the earliest chemical abundances imprinted in the oldest stars, and the nature of the supernovae that enriched them. Most of the current observational efforts in finding the chemical imprints left by the first stars have focused on the most metal-poor stars found in the Milky Way halo (Beers \& Christlieb 2005; Beers et al. 2017). Very metal-poor stars were also found in ultra-faint dwarf galaxies, which are intriguing dark-matter dominated objects with very low average metallicities (Ji et al. 2016). The Galactic bulge, as well as the halo, is a potential host of some of the oldest stars in our Galaxy. Tumlinson (2010) suggests that half of the oldest stars were formed in the central parts of the Galaxy. Searches for field metalpoor stars in the Galactic bulge are the target of surveys such as those by Howes et al. (2016), Casey \& Schlaufman (2015), the Pristine Inner Galaxy Survey (PIGS, Arentsen et al. 2020), HERBS (Duong et al. 2019a,b), and COMBS (Lucey et al. 2019, 2021, 2022) surveys. Metal-poor stars in the Galactic bulge have been mostly traced by Globular clusters (Rossi et al. 2015, Bica et al. 2016), and RR Lyrae stars (Minniti et al. 2017), which show a peak at $[\mathrm{Fe} / \mathrm{H}] \sim-1.0$ (Barbuy et al. 2018a). This metallicity peak at $[\mathrm{Fe} / \mathrm{H}] \sim-1.0$ has been also recently confirmed regarding field stars by Lucey et al. (2021). In fact, it is expected that a fast chemical enrichment in the Galactic bulge results in a very old population with this relatively high metallicity, that would correspond to the age of stars with $[\mathrm{Fe} / \mathrm{H}] \sim-3.0$ in the halo (Chiappini et al. 2011, Wise et al. 2012, Barbuy et al. 2018a).

Our main interest in the present work is to analyse the abundances of stars of the spheroidal bulge with a moderate metallicity of $[\mathrm{Fe} / \mathrm{H}]<-0.8$, in order to try to identify the earliest supernovae of the central regions of the Galaxy, and imposing constraints on the early chemical enrichment of the Milky Way. For the selection of sample stars we applied kinematical and dynamical criteria, by combining data from APOGEE and Gaia Early Release EDR3. We
chose stars with azimuthal velocity $V_{\phi}<0$ (this selection will avoid contamination by disk stars, but would still include accreted debris of objects such as GES) that have orbits confined within 4 kpc of the Galactic center, a maximum height of $|\mathrm{z}|_{\max }<3.0 \mathrm{kpc}$, eccentricity $>0.7$, and with orbits not supporting the bar structure. With this selection, as noted above, we expect our sample to be dominated by a pressure supported, most probably old component of the bulge. We hope to discard the contamination of our sample by accreted debris thanks to the detailed chemical information and, in particular, the alpha-over-iron enhancement, expected to be low in most of the accreted debris. Finally, given that we used a barred potential, the z-component of angular momentum (Lz) is not conserved, and most orbits are either retrograde or prograde, and a fraction among those identified as counter-rotating keep retrograde along its orbit.

In this paper we carried out an analysis of atomic and molecular lines for the selected sample of 58 metal-poor spheroid bulge star candidates aiming at refining the APOGEE Stellar Parameter and Chemical Abundance Pipeline (ASPCAP; García-Pérez et al. 2016) results, in order to interpret the derived abundances in terms of the early chemo-dynamical evolution of the bulge. As it will be shown, this re-analysis is critical for some alpha elements, and therefore for the identification and confirmation of old spheroid bulge stars at moderately low metallicities. In the present work we adopt the stellar parameters issued from the DR17 release of the APOGEE ASPCAP code. The $\mathrm{C}, \mathrm{N}$, and O abundances are derived from $\mathrm{CO}, \mathrm{OH}$ and CN lines, that are interdependent, and since there are such molecular lines all over the spectra, they can affect the abundances of atomic lines. We also refine the abundances of Ce . Other elements including $\mathrm{Na}, \mathrm{Al}$ and iron-peak elements will be the topic of a future work.

In Section 2 the selection of our sample is described. The element abundances are derived in Sect. 3. In Sect. 4 the results are compared with literature data for bulge samples and chemodynamical models, and discussed. In Sect. 5 conclusions are drawn.

## 2 THE SAMPLE

The Apache Point Observatory Galactic Evolution Experiment (APOGEE; Majewski et al. 2017) is part of the Sloan Digital Sky Survey III and IV (SDSS; Blanton et al. 2017). It is a project encompassing spectroscopic programs that observe Milky Way stars at high resolution and high signal-to-noise ratios $(\mathrm{S} / \mathrm{N})$ in the near-infrared (NIR). The project SDSS-IV technical summary, the SDSS telescope and APOGEE spectrograph are described in Blanton et al. (2017), Gunn et al. (2006), and Wilson et al. (2019), respectively, whereas Zasowski et al. (2013, 2017), Beaton et al. (2020) and Santana et al. (2021) describe the APOGEE and APOGEE-2 Target Selections. The data release 17 (DR17) contains high-resolution ( $\mathrm{R} \sim 22,500$ ) NIR spectra ( $15140-16940 \AA$ ) for some $7 \times 10^{5}$ stars, covering both the northern and southern sky. While APOGEE-1 observed the Milky Way bulge/bar at $l>0 \mathrm{deg}$, APOGEE- 2 covers the whole bulge/bar region.

Given that the central part of the Milky Way hosts members of all Galactic components, including the bulge, disc, and halo (PérezVillegas et al. 2020; Rojas-Arriagada et al. 2020; Queiroz et al. 2021), we have used the chemo-orbital analysis shown in Queiroz et al. (2021) to identify good candidates in the spheroidal bulge APOGEE sample. To disentangle the different stellar populations coexisting in the innermost parts of the Galaxy is not an easy task, and one of the difficulties is to compute precise distances for these stars due to the high extinction. Thanks to StarHorse (Santiago et al. 2016; Queiroz et al. 2018), precise stellar distances for the entire APOGEE sample
were derived both for DR16 (Ahumada et al. 2020, Queiroz et al. 2020), and DR17 ${ }^{1}$ (Abdurro'uf et al. 2022, Queiroz et al. 2022, in prep.).

We selected stars from the reduced-proper-motion (RPM) sample of Queiroz et al. (2021). For that sample, orbits were calculated using the StarHorse distances and the proper motions from the Gaia Early Data Release 3 (EDR3) (Gaia Collaboration 2021). In order to select the best candidate objects that belong to the spheroidal bulge, the following selection criteria were adopted: a maximum distance to the Galactic center of $\mathrm{d}_{\mathrm{GC}}<4 \mathrm{kpc}$ (Bica et al. 2016); a maximum vertical excursion from the Galactic plane $|z|_{\text {max }}<3.0 \mathrm{kpc}$; eccentricity $>0.7$; orbits that do not support the bar structure ${ }^{2}$ (orbits with frequency ratio $f_{R} / f_{x} \neq 2.0 \pm 0.1$; Portail et al. 2015); and based on Figure 17 of Queiroz et al. (2021), we selected counterrotating stars ( $V_{\phi}<0.0$ ). Finally, according to the discussion of Sect. 1, we considered only stars with moderate metallicity of $[\mathrm{Fe} / \mathrm{H}]$ $<-0.80$. Applying the selection criteria described above, a sample of 58 stars has been selected. The adopted input parameters for the orbits integration and the orbital parameters are given in Table 1. In Figure 1 we show the distribution of parameters for our selected stars in comparison with the RPM sample of Queiroz et al. (2021) and our selection is then similar to the metal-poor/high eccentricity stars discussed in their Fig. 20. This figure indicates that our selection is indeed reaching bulge stars of the metal-poor spheroid, that are moderately metal-poor, $\alpha$-rich and in eccentrical orbits but confined to the Galactic center region. Figure 2 shows the projected $1, b$ distribution of the sample in the Galactic bulge region.

As explained above, our stars were selected from the reduced-proper-motion sample of Queiroz et al. (2021), and therefore have a signal-to-noise $S N R>50$, a good spectral fit from the ASPCAP pipeline ASPCAP_Chi $2<25$, and a radial velocity scatter Vscatter $<1.5 \mathrm{~km} \mathrm{~s}^{-1}$. As for the renormalized unit weight error RUWE Gaia EDR3 parameter, 56 out of 58 stars in our sample comply with the standard or minimal requirements to get reliable orbital elements, since astrometry from Gaia EDR3 has its own caveats. According to the Gaia consortium, the RUWE parameter is suggested to return stars astrometrically well-behaved by applying a cut with RUWE $\leqslant 1.4$, which is followed by the 56 stars listed in Table 1. The stars 2M17453659-2309130 and 2M18023156-2834451 have a RUWE $>1.4$, which makes them sources with astrometric parameters that are not reliable enough.

In Figure 3 a Kiel diagram of the sample stars is plotted with the effective temperature from Apogee-ASPCAP and gravity log g coming from the StarHorse output from Queiroz et al. (2020), and compared with the reduced-proper-motion sample of Queiroz et al. (2021).

## 3 ANALYSIS

We have initially adopted the calibrated stellar parameters effective temperature $\mathrm{T}_{\text {eff }}$, gravity $\log \mathrm{g}$, metallicity $[\mathrm{Fe} / \mathrm{H}]$ and microturbulence velocity $\mathrm{v}_{t}$ from APOGEE DR16 - we point out that the calibrated parameters give very different element abundances, and should not be used for such aims. In fact the results from the APOGEE Stellar Parameters and Chemical Abundances Pipeline (ASPCAP) (GarcíaPérez et al. 2016) are obtained for the reported non-calibrated spectroscopic stellar parameters. We then adopted these non-calibrated

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Figure 1. Comparison of the present sample of 58 selected stars (red) and the RPM sample of Queiroz et al. (2021) (blue). Upper panels: normalized distribution of metallicity and alpha-to-iron ratios from APOGEE; lower panels: mean radius $\mathrm{R}_{\text {mean }}=\left(\left(\mathrm{R}_{\text {apocenter }}+\mathrm{R}_{\text {pericenter }}\right) / 2\right.$ and eccentricity of the orbits.


Figure 2. Projected 1,b distribution of studied stars in the Galactic bulge region. Symbols: filled stars: this work; filled circles: bulge globular clusters (GCs); solid black line: contours of the bulge. The colours indicate metallicity according to the colour-bar.
stellar parameterers from DR17, since we became aware that these are obtained from a spectroscopic solution that minimizes the errors in 7 dimensions ( $\left.\mathrm{T}_{\text {eff }}, \log \mathrm{g},[\mathrm{Fe} / \mathrm{H}], \mathrm{v}_{t},[\alpha / \mathrm{Fe}],[\mathrm{C} / \mathrm{Fe}],[\mathrm{N} / \mathrm{Fe}]\right)$.

For this reason we proceeded with all the rederivation of abundances with the DR17 non-calibrated parameters. These stellar parameters are reported in Table 2, and they are the final parameters adopted.

The abundances were determined by comparing the observed spectra with the synthetic ones. The synthetic spectra calculations are carried out with the code PFANT ${ }^{3}$, as described in Barbuy et al. (2018b). This code is an update of the original FANTOM or ABON2 Meudon code by M. Spite. Each model atmosphere was interpolated in the MARCS grids (Gustafsson et al. 2008).

The atomic line list employed is that from the APOGEE collaboration (Smith et al. 2021). Molecular electronic transition lines of CN $\mathrm{A}^{2} \Pi-\mathrm{X}^{2} \Sigma$, vibration-rotation $\mathrm{CO}^{1} \Sigma^{+}, \mathrm{OH} \mathrm{X}^{2} \Pi$ and $\mathrm{TiO} \phi$-system

[^3]Table 1. Coordinates, Starshorse distances, Gaia EDR3 proper motions, Gaia DR2 radial velocity, and orbital parameters for the selected 58 stars from RPM sample of Queiroz et al. (2021).

| ID | $\begin{gathered} \alpha \\ \left({ }^{\circ}\right) \end{gathered}$ | $\begin{gathered} \delta \\ \left({ }^{\circ}\right) \end{gathered}$ | $\underset{(\mathrm{kpc})}{\mathrm{d}_{\odot}}$ | $\left(\text { mas yr }^{-1}\right)$ | $\begin{gathered} \mu_{\delta} \\ \left(\text { mas yr}^{-1}\right) \end{gathered}$ | $\begin{gathered} \mathrm{RV} \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} r_{\min } \\ (\mathrm{kpc}) \end{gathered}$ | $\begin{gathered} r_{\max } \\ (\mathrm{kpc}) \end{gathered}$ | $\begin{gathered} \|z\|_{\max } \\ (\mathrm{kpc}) \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2M17153858-2759467 | 258.911 | -27 | $8.51 \pm 0.50$ | $-5.46 \pm 0.02$ | $-5.30 \pm 0.02$ | $191.79 \pm 0.01$ | $0.13 \pm 0.05$ | $2.51 \pm 0.46$ | $1.68 \pm 0.18$ | $0.90 \pm 0.04$ |
| 2 M | 259.385 | -25.315 | 91 | $-2.14 \pm 0.04$ | $-9.47 \pm 0.03$ | $187.54 \pm 0.02$ | $0.19 \pm 0.16$ | 79 | 49 | 05 |
| 2M17173693-2806495 | 259.404 | -28.114 | $6.94 \pm 0.45$ | $-4.85 \pm 0.03$ | $-9.80 \pm 0.02$ | -104.63 | $0.12 \pm 0.06$ | $1.94 \pm 0.44$ | $1.52 \pm 0.19$ | $0.89 \pm 0.05$ |
| 2M17190320-2857321 | 259.763 | -28.959 | $6.81 \pm 0.46$ | $-5.95 \pm 0.03$ | $-7.60 \pm 0.02$ | $-83.87 \pm 0.03$ | $0.14 \pm 0.07$ | $1.80 \pm 0.33$ | $0.74 \pm 0.26$ | $0.87 \pm 0.05$ |
| 2M17224443-2343053 | 260.685 | -2 | $6.02 \pm 0.42$ | $-9.20 \pm 0.02$ | $-8.15 \pm 0.01$ | $114.23 \pm 0.01$ | 12 | $3.87 \pm 0.51$ | $2.61 \pm 0.23$ | . 77 |
| 2 M | 261.262 | -28.011 | $5.83 \pm 0.76$ | $-3.05 \pm 0.03$ | -9. | 26. | $0.20 \pm 0.11$ | $2.63 \pm 0.71$ | $1.02 \pm 0.38$ | 07 |
| 2M17265563-2813558 | 261.732 | -28.232 | $7.55 \pm 0.56$ | $-7.25 \pm 0.04$ | $-7.31 \pm 0.03$ | $196.52 \pm 0.03$ | $0.13 \pm 0.07$ | $2.40 \pm 0.53$ | 36 | $0.91 \pm 0.04$ |
| 2M17281191-2831393 | 262.050 | -28.528 | $6.50 \pm 0.58$ | $-9.70 \pm 0.03$ | $-4.61 \pm 0.02$ | $81.01 \pm 0.02$ | $0.14 \pm 0.05$ | $2.24 \pm 0.68$ | $1.87 \pm 0.29$ | $0.90 \pm 0.05$ |
| 2M17285088-2855427 | 262.212 | -28.929 | $7.59 \pm 0.42$ | $-4.80 \pm 0.03$ | $-5.57 \pm 0.02$ | $-7.43 \pm 0.01$ | $0.06 \pm 0.03$ | $0.83 \pm 0.25$ | $0.47 \pm 0.02$ | $0.87 \pm 0.04$ |
| 2M17291778-2602468 | 262.324 | -26.046 | $6.93 \pm 0.44$ | $-5.60 \pm 0.06$ | $-7.06 \pm 0.04$ | $-47.65 \pm 0.01$ | $0.12 \pm 0.07$ | $1.50 \pm 0.43$ | . 10 | . 7 |
| 2M17292082-2126433 | 262.337 | -21.445 | $6.60 \pm 0.68$ | $-0.84 \pm 0.02$ | $-10.79 \pm 0.02$ | $-79.08 \pm 0.01$ | $0.19 \pm 0.10$ | $2.83 \pm 0.67$ | $2.06 \pm 0.14$ | $0.87 \pm 0.09$ |
| 2M | 262.395 | -27.688 | $6.81 \pm 0.52$ | $-3.56 \pm 0.04$ | $-8.16 \pm 0.03$ | $-74.26 \pm 0.02$ | $0.09 \pm 0.06$ | $1.51 \pm 0.39$ | $0.59 \pm 0.10$ | $0.89 \pm 0.05$ |
| 2M17295481-2051262 | 262.478 | -20.857 | $7.00 \pm 0.38$ | $0.11 \pm 0.04$ | $-6.20 \pm 0.03$ | $-213.15$ | $0.16 \pm 0.08$ | $3.43 \pm 0.40$ | $2.31 \pm 0.08$ | $0.90 \pm 0.04$ |
| 2M17301495-2337002 | 262.562 | -23.617 | $8.28 \pm 0.66$ | $-8.24 \pm 0.04$ | $-9.11 \pm 0.02$ | $-70.19 \pm 0.01$ | $0.24 \pm 0.17$ | $1.97 \pm 0.93$ | $1.81 \pm 0.32$ | $0.83 \pm 0.17$ |
| 2M17303581-2354453 | 262.649 | -23.913 | $7.99 \pm 0.60$ | $-8.31 \pm 0.04$ | $-4.45 \pm 0.02$ | $27.88 \pm 0.01$ | $0.10 \pm 0.14$ | $1.52 \pm 0.40$ | $1.38 \pm 0.18$ | 08 |
| 2M17310874-2956542 | 262.786 | -29.948 | $6.81 \pm 0.00$ | $-3.38 \pm 0.04$ | $-7.93 \pm 0.03$ | $-10.11 \pm 0.02$ | $0.19 \pm 0.01$ | $1.50 \pm 0.01$ | $0.36 \pm 0.00$ | $0.77 \pm 0.01$ |
| 2M17323787-2023013 | 263.158 | -20.384 | $7.76 \pm 0.55$ | $-5.22 \pm 0.03$ | $-1.34 \pm 0.02$ | $-97.24 \pm 0.01$ | $0.17 \pm 0.08$ | $2.53 \pm 0.24$ | $1.62 \pm 0.22$ | $0.88 \pm 0.06$ |
| 2M17324257-2301417 | 263.177 | -23.028 | $7.69 \pm 0.74$ | $-2.70 \pm 0.05$ | $-7.92 \pm 0.03$ | $-181.81 \pm 0.01$ | $0.16 \pm 0.13$ | $1.67 \pm 0.58$ | 0.28 | . 09 |
| 2M17330695-2302130 | 263.279 | -23.037 | $7.40 \pm 0.10$ | $-3.51 \pm 0.04$ | $-9.38 \pm 0.03$ | $6.42 \pm 0.00$ | $0.11 \pm 0.05$ | $1.44 \pm 0.04$ | . $95 \pm 0.01$ | $0.86 \pm 0.05$ |
| 2M17330730-2407378 | 263.280 | -24.127 | $5.32 \pm 0.25$ | $-4.74 \pm 0.03$ | $-8.85 \pm 0.02$ | $-31.23 \pm 0.01$ | $0.11 \pm 0.04$ | $3.15 \pm 0.25$ | $1.41 \pm 0.42$ | . $93 \pm 0.03$ |
| 2 M | 263 | -39.086 | $8.63 \pm 0.69$ | $-2.19 \pm 0.07$ | $-3.37 \pm 0.05$ | . 03 | $0.23 \pm 0.06$ | $1.90 \pm 0.22$ | $0.65 \pm 0.27$ | $0.78 \pm 0.05$ |
| 2M17342067-3902066 | 263.586 | -39.035 | $9.80 \pm 0.00$ | $-2.51 \pm 0.08$ | $-3.17 \pm 0.06$ | $5.95 \pm 0.04$ | $0.13 \pm 0.05$ | $2.50 \pm 0.14$ | $1.50 \pm 0.18$ | $0.90 \pm 0.04$ |
| 2M17344841-4540171 | 263.702 | -45.671 | $6.71 \pm 0.38$ | $-0.85 \pm 0.02$ | $-6.51 \pm 0.01$ | $148.00 \pm 0.01$ | $0.17 \pm 0.15$ | $3.67 \pm 0.51$ | 18 | . 09 |
| 2M17351981-1948329 | 263.833 | -19.809 | $8.20 \pm 0.32$ | $-2.39 \pm 0.02$ | $-6.57 \pm 0.01$ | $-230.13 \pm 0.00$ | $0.36 \pm 0.23$ | $2.61 \pm 0.53$ | $2.14 \pm 0.16$ | $0.77 \pm 0.16$ |
| 2M17354093-1716200 | 263.921 | -17.272 | $6.15 \pm 0.35$ | $-4.18 \pm 0.02$ | $-7.53 \pm 0.01$ | $-84.29 \pm 0.01$ | $0.17 \pm 0.11$ | $2.84 \pm 0.27$ | $1.59 \pm 0.16$ | $0.88 \pm 0.06$ |
| 2 M | 264. | -24.405 | $6.78 \pm 0.52$ | $-2.34 \pm 0.07$ | $-8.58 \pm 0.04$ | $-56.51 \pm 0.01$ | $0.12 \pm 0.09$ | $1.68 \pm 0.43$ | 0.08 | 10 |
| 2M17390801-2331379 | 264.783 | -23.527 | $7.57 \pm 0.54$ | $-7.05 \pm 0.03$ | $-3.91 \pm 0.02$ | $-199.67 \pm 0.01$ | $0.13 \pm 0.07$ | $1.92 \pm 0.36$ | $1.43 \pm 0.17$ | $0.88 \pm 0.05$ |
| 2M1 | 264.863 | -23.175 | $6.70 \pm 0.31$ | $-10.26 \pm 0.03$ | $-7.39 \pm 0.02$ | $47.66 \pm 0.00$ | $0.13 \pm 0.07$ | $2.67 \pm 0.47$ | 22 | . 04 |
| 2M17453659-2309130 | 266.402 | -23.154 | $6.31 \pm 0.56$ | $-4.98 \pm 0.23$ | $-7.39 \pm 0.15$ | $-140.43 \pm 0.02$ | $0.12 \pm 0.04$ | $2.21 \pm 0.47$ | $0.54 \pm 0.32$ | $0.90 \pm 0.05$ |
| 2M17473299-2258254 | 266.887 | -22.974 | $7.36 \pm 0.61$ | $-4.18 \pm 0.02$ | $-9.24 \pm 0.01$ | $-39.26 \pm 0.01$ | $0.11 \pm 0.04$ | $1.37 \pm 0.35$ | $0.44 \pm 0.04$ | $0.87 \pm 0.05$ |
| 2M17482995-2305299 | 267.125 | -23.092 | $7.05 \pm 0.43$ | $-0.95 \pm 0.03$ | $-6.72 \pm 0.02$ | -216 | 0.1 | $2.07 \pm 0.51$ | . 32 | . 06 |
| 2M17483633-2242483 | 267.151 | -22.713 | $8.11 \pm 0.69$ | $-0.62 \pm 0.03$ | $-9.74 \pm 0.02$ | $-93.04 \pm 0.00$ | $0.12 \pm 0.07$ | $1.38 \pm 0.61$ | $0.87 \pm 0.34$ | $0.85 \pm 0.09$ |
| 2M | 267.628 | -23.223 | $6.83 \pm 0.38$ | $-4.88 \pm 0.05$ | $-6.57 \pm 0.03$ | $-203.16 \pm 0.01$ | $0.09 \pm 0.03$ | $1.94 \pm 0.42$ | $37 \pm 0.13$ | 03 |
| 2M17503263-3654102 | 267.636 | -36.903 | $7.49 \pm 0.62$ | $-7.00 \pm 0.02$ | $-4.97 \pm 0.01$ | $11.58 \pm 0.01$ | $0.10 \pm 0.05$ | $1.51 \pm 0.38$ | $1.28 \pm 0.08$ | $0.89 \pm 0.04$ |
| 2M17511568-3249403 | 267.815 | -32.828 | $7.54 \pm 0.59$ | $-4.58 \pm 0.04$ | $-9.25 \pm 0.03$ | $-102.21 \pm 0.01$ | $0.08 \pm 0.03$ | $1.28 \pm 0.21$ | $0.45 \pm 0.03$ | $0.88 \pm 0.04$ |
| 2M17532599-2053 | 26 | -20.892 | 7. | $-3.44 \pm 0.04$ | $-7.77 \pm 0.03$ | $-78.10 \pm 0.01$ | 0.08 | 1.3 | 0.05 | $0.89 \pm 0.04$ |
| 2M17552681-3334272 | 268.862 | -33.574 | $7.67 \pm 0.55$ | $-3.57 \pm 0.03$ | $-4.88 \pm 0.02$ | $166.48 \pm 0.02$ | $0.09 \pm 0.03$ | $1.43 \pm 0.25$ | $0.65 \pm 0.04$ | $0.89 \pm 0.04$ |
| 2 M | 268 | -32 | $7.10 \pm 0.89$ | $-7.00 \pm 0.03$ | $-6.81 \pm 0.02$ | -7 | 05 | $1.34 \pm 0.68$ | 24 | . 06 |
| 2M18005152-2916576 | 270.215 | -29.283 | $8.45 \pm 0.60$ | $1.18 \pm 0.04$ | $-9.34 \pm 0.03$ | $-77.43 \pm 0.02$ | $0.17 \pm 0.09$ | $1.25 \pm 0.61$ | $1.07 \pm 0.34$ | $0.80 \pm 0.10$ |
| 2M18010424-3126158 | 270.268 | -31.438 | $7.10 \pm 0.57$ | $-1.22 \pm 0.03$ | $-9.10 \pm 0.02$ | $81.96 \pm 0.00$ | $0.09 \pm 0.04$ | $1.43 \pm 0.50$ | $0.82 \pm 0.15$ | $0.88 \pm 0.06$ |
| 2M18020063-1814495 | 270.503 | -18.247 | $5.97 \pm 0.38$ | $-4.65 \pm 0.05$ | $-8.19 \pm 0.04$ | $-94.07 \pm 0.02$ | $0.11 \pm 0.03$ | $2.85 \pm 0.40$ | 55 | $0.92 \pm 0.02$ |
| 2M18023156-2834451 | 270.632 | -28.579 | $8.15 \pm 0.44$ | $-4.55 \pm 0.07$ | $-10.32 \pm 0.05$ | $-190.17 \pm 0.01$ | $0.27 \pm 0.10$ | $1.59 \pm 0.40$ | $0.62 \pm 0.33$ | $0.73 \pm 0.12$ |
| 2M18042687-2928348 | 271.112 | -29.476 | $7.89 \pm 0.75$ | $-2.34 \pm 0.03$ | $-7.82 \pm 0.02$ | $-113.51 \pm 0.02$ | $0.08 \pm 0.06$ | $1.05 \pm 0.41$ | $0.61 \pm 0.07$ | $0.89 \pm 0.07$ |
| 2M18044663-3132174 | 271.194 | -31.538 | $7.31 \pm 0.43$ | $-6.68 \pm 0.03$ | $-7.25 \pm 0.02$ | $-145.20 \pm 0.01$ | $0.10 \pm 0.04$ | $1.62 \pm 0.30$ | $1.10 \pm 0.14$ | $0.89 \pm 0.05$ |
| 2M18050452-3249149 | 271.269 | -32.821 | $5.51 \pm 0.41$ | $-3.19 \pm 0.02$ | $-10.36 \pm 0.01$ | $46.90 \pm 0.01$ | $0.12 \pm 0.06$ | $3.58 \pm 0.59$ | $1.47 \pm 0.51$ | $0.93 \pm 0.04$ |
| 2M18050663-3005419 | 271.278 | -30.095 | $7.92 \pm 0.36$ | $-1.98 \pm 0.04$ | $-8.42 \pm 0.03$ | $-137.47 \pm 0.00$ | $0.09 \pm 0.04$ | $1.09 \pm 0.11$ | $0.77 \pm 0.05$ | $0.86 \pm 0.07$ |
| 2M18052388-2953056 | 271.350 | -29.885 | $7.43 \pm 0.59$ | $-5.77 \pm 0.03$ | $-8.14 \pm 0.02$ | $-4.77 \pm 0.05$ | $0.09 \pm 0.06$ | $1.11 \pm 0.26$ | $0.66 \pm 0.03$ | $0.85 \pm 0.06$ |
| 2M18065321-2524392 | 271.722 | -25.411 | $7.91 \pm 0.80$ | $-7.64 \pm 0.06$ | $-8.61 \pm 0.04$ | $-112.08 \pm 0.01$ | $0.28 \pm 0.14$ | $1.71 \pm 0.83$ | $0.60 \pm 0.46$ | $0.74 \pm 0.13$ |
| 2M18080306-3125381 | 272.013 | -31.427 | $10.06 \pm 0.73$ | $-1.89 \pm 0.05$ | $-4.50 \pm 0.04$ | $23.35 \pm 0.03$ | $0.13 \pm 0.08$ | $2.39 \pm 0.80$ | $1.38 \pm 0.43$ | $0.89 \pm 0.05$ |
| 2M18104496-2719514 | 272.687 | -27.331 | $7.30 \pm 0.31$ | $-1.79 \pm 0.03$ | $-7.09 \pm 0.03$ | $-163.54 \pm 0.02$ | $0.12 \pm 0.05$ | $1.57 \pm 0.27$ | $0.63 \pm 0.05$ | $0.87 \pm 0.05$ |
| 2M18125718-2732215 | 273.238 | -27.539 | $8.12 \pm 0.34$ | $-5.57 \pm 0.02$ | $-7.99 \pm 0.02$ | $-86.39 \pm 0.00$ | $0.13 \pm 0.04$ | $1.14 \pm 0.16$ | $0.79 \pm 0.09$ | $0.80 \pm 0.07$ |
| 2M18142265-0904155 | 273.594 | -9.071 | $6.92 \pm 0.26$ | $-1.42 \pm 0.11$ | $-8.72 \pm 0.09$ | $-151.19 \pm 0.02$ | $0.28 \pm 0.13$ | $3.68 \pm 0.28$ | $2.50 \pm 0.59$ | $0.87 \pm 0.06$ |
| 2M18143710-2650147 | 273.655 | -26.837 | $7.46 \pm 0.52$ | $-3.66 \pm 0.04$ | $-7.43 \pm 0.04$ | $-200.72 \pm 0.02$ | $0.16 \pm 0.13$ | $1.62 \pm 0.38$ | $1.07 \pm 0.25$ | $0.85 \pm 0.14$ |
| 2M18150516-2708486 | 273.772 | -27.147 | $6.80 \pm 0.38$ | $-0.00 \pm 0.03$ | $-9.08 \pm 0.02$ | $-141.63 \pm 0.02$ | $0.13 \pm 0.05$ | $2.28 \pm 0.35$ | $1.25 \pm 0.22$ | $0.88 \pm 0.04$ |
| 2M18195859-1912513 | 274.994 | -19.214 | $6.07 \pm 0.29$ | $-6.28 \pm 0.05$ | $-6.91 \pm 0.03$ | $-78.88 \pm 0.02$ | $0.10 \pm 0.05$ | $2.81 \pm 0.24$ | $0.76 \pm 0.37$ | $0.93 \pm 0.03$ |
| 2M18200365-3224168 | 275.015 | -32.405 | $6.27 \pm 0.40$ | $-4.20 \pm 0.03$ | $-10.36 \pm 0.02$ | $-124.55 \pm 0.01$ | $0.18 \pm 0.08$ | $3.35 \pm 0.47$ | $2.19 \pm 0.21$ | $0.91 \pm 0.05$ |
| 2M18344461-2415140 | 278.686 | -24.254 | $7.46 \pm 0.50$ | $-3.96 \pm 0.03$ | $-8.27 \pm 0.02$ | $-171.66 \pm 0.02$ | $0.23 \pm 0.20$ | $2.63 \pm 0.49$ | $2.12 \pm 0.12$ | $0.85 \pm 0.14$ |
| 2M18500307-1427291 | 282.513 | -14.458 | $6.40 \pm 0.32$ | $0.13 \pm 0.03$ | $-6.71 \pm 0.02$ | $-134.03 \pm 0.02$ | $0.16 \pm 0.07$ | $3.92 \pm 0.29$ | $2.49 \pm 0.28$ | $0.92 \pm 0.03$ |

Table 2. Selected 58 stars and corresponding DR17 non-calibrated stellar parameters.

| ID | $\mathrm{T}_{\text {eff(nc) }}$ | $\log \mathrm{g}_{(\mathrm{nc})}$ | $[\mathrm{Fe} / \mathrm{H}]_{(\mathrm{nc})}$ | $\mathrm{v}_{t}$ |
| :--- | :---: | :---: | :---: | :---: |
|  | $(\mathrm{~K})$ |  |  | $(\mathrm{km} / \mathrm{s})$ |



Figure 3. Kiel diagram of the 58 sample bulge stars. (purple circles). In the background, we show the full reduced proper-motion sample of Queiroz et al. (2021).
$\mathrm{b}^{1} \Pi-\mathrm{d}^{1} \Sigma$ lines were included. The line lists for CN were made available by S. P. Davis, the CO line lists were adopted from Goorvitch (1994), and the OH are from Goldman et al. (1998). For TiO the line list is from Jorgensen (1994). More details on CN, CO, OH and TiO molecular lines are given in Meléndez \& Barbuy (1999), Meléndez et al. (2001, 2002, 2003), Schiavon \& Barbuy (1999) and Barbuy et al. (2018b).

The atomic lines analysed initially were selected from Smith et al. (2021), Shetrone et al. (2015), Ce II lines identified by Cunha et al. (2017), and lines of S I identified by Fanelli et al. (2021). Lines of Nd II (Hasselquist et al. 2016) and Yb II (Smith et al. 2021) were not studied. In Table 3 are reported the lines that we verified in the spectra of the 58 sample stars.

For the moderately metal-poor sample stars, some of the lines indicated in the articles above are not suitable, and in a few cases we have added other lines that we identified as suitable for the stellar parameters of the sample stars. The lines are discussed in detail below. In the present work we adopt the ASPCAP abundances of $\mathrm{Mg}, \mathrm{Si}$, Ca and revise the $\mathrm{C}, \mathrm{N}, \mathrm{O}$, and Ce abundances; we also verified Ti lines and some comments are given, but the abundances are not used, given conflicting results from different lines. Other elements such as $\mathrm{Na}, \mathrm{Al}$ and iron-peak elements will be analysed elsewhere.

We identified and fitted the studied lines in the reference stars Arcturus and $\mu$ Leo, in order to check if the lines are well reproduced in these stars, and therefore reliable for deriving abundances in the sample stars. For the reference star Arcturus, we used the Hinkle et al. (1995) atlas, and for the metal-rich reference giant star $\mu$ Leo a spectrum from APOGEE was used. The adopted stellar parameters for Arcturus and $\mu$ Leo are from Meléndez et al. (2003) and Zoccali et al. (2006) plus Lecureur et al. (2007), respectively.

Table 4 reports abundances in the Sun, Arcturus and $\mu$ Leo. For the Sun they are from a) Grevesse et al. 1996, 1998, adopted, b) Grevesse et al. (2015), Scott et al. (2015a,b), c) Lodders et al. (2009). For Arcturus, the abundances are from Meléndez et al. (2003), Lecureur
et al. (2007), Ramírez \& Allende Prieto (2011), Barbuy et al. (2014), and Smith et al. (2013). For $\mu$ Leo, the abundances are from, Gratton \& Sneden (1990), Smith \& Ruck (2000), Lecureur et al. (2007), Barbuy et al. (2015), Smith et al. (2013) or present fits, using the observed spectrum by Lecureur et al. (2007) in the optical.

According to Ashok et al. (2021), and Nidever et al. (2015) the average resolution of the APOGEE observations is $R \approx 22,500$ based on a direct-measured FWHM of $\sim 0.7 \AA$, with $10-20 \%$ variations seen across the wavelength range. We have employed a typical FWHM $=0.70 \AA$, but to fit better different lines we varied the FWHM values from 0.60 to 0.75 , from the lowest to the highest wavelengths. Note that the FWHM varies from fibers to fibers and with a fiber with wavelength.

### 3.1 C, N, O abundances

The abundances of $\mathrm{C}, \mathrm{N}$ and O are derived from $\mathrm{CN}, \mathrm{OH}$ and CO molecular lines. They are interdependent due to the molecular dissociative equilibrium. Since the molecular lines are spread all over the spectra, these abundances are derived first, and they are reported in Table 5.

Computing synthetic spectra employing the PFANT code described in Barbuy et al. (2018b), we have derived C, N, O abundances in two ways:
method $a$ ): in the region 15144-16896 $\AA$, first we derive the O abundances by analysing the molecular lines of OH . Some of the most prominent OH lines in this region are at: $15264.60,15266.160,15278.516,15281.045,15719.687$, $15893.524,16074.151,16662.187,16872.265,16895.164$ Å. These lines are the most sensitive to oxygen variation in the APOGEE sample. We derive C abundances by analysing the CO molecular lines, but there are not many strong CO lines in the range of 15100-17000 $\AA$ A. In our sample, the strongest lines of CO, used to measure C abundances, are at $15983.214,15985.598,15990.420,16016.081$ Å. Next we see how the CN lines change when we modify Nitrogen. The most sensitive lines of CN are at $15162.648,15222.382 \AA$. There are many CN lines in the region we are working with (especially in the range 15522-15600 $\AA$ ), but most of them are too shallow to give reliable abundance measurements.
method $b$ ): a derivation of CNO abundances using the region 15525-15595 $\AA$, where there are clear lines of OH , and a clear bandhead of CO, as well as lines of CN , although less conspicuous, as done for example in Barbuy et al. (2021a) for Phoenix spectra that were observed in this region only. For these calculations a FWHM $=0.60$ was adopted which is suitable for the wavelength region in question.

This is illustrated in Figure 4 for star 2M17382504-2424163. Note the clear OH lines at 15535.46, 15565.91, 15566.78, and clear CO bandhead at $15577.4 \AA$.

We concluded that both methods a) and b) give very similar results within $\pm 0$. 1 dex.

A verification of these CNO abundances was carried out by fitting lines along all the spectra, in particular the lines of CO $15600.74,15612.5,15667.55 \AA$, where only for four stars the C abundance was decreased by -0.05 to -0.10 (stars 2M17330695-2302130, 2M18050663-3005419, 2M18125718-2732215 and 2M183444612415140), and for the others the fits were very satisfactory.

We then proceeded with the verification of the OH lines: OH 15130.921, 15266.168, 15281.052, 15409.172, 15568.78, 15651.896, 15719.696, 15755.522 A, and CN 5181.277, 15298.487, 15308.893, 15318.74, 15337.959, 15341.508, 15432.811, $15447.095,15466.235,15481.868,15530.776,15684.088$, 15737.445 A. Only for star 2M18023156-2834451 we increased the
oxygen abundance by 0.05 dex, noting that its spectra shows larger lines than the others, needing a higher spectral convolution to be fitted.

Fits are shown for selected OH lines for star 2M17382504-2424163 in Figure 5, and CO lines in Figure 6 for stars 2M17382504-2424163 and 2M17511568-3249403.

Regions of CN lines are verified, using wavelength regions indicated by Fernández-Trincado et al. (2020a,b) for example. In Figure 7 are shown the fits to good CN lines. Among these, the clearest CN feature is at $15387.6 \AA$, and its fits are compatible with the C,N abundances from the 15283-15287, 15320-15330 and 15355-15380 $\AA$ regions. The feature at $15514 \AA$ is blended with a CoI line and is less reliable. The N abundances derived are confirmed for about half the stars, whereas for the other half the N abundance was decreased by a mean of 0.2 dex : this is not surprising because the CN lines in the $15555 \pm 50 \AA$ from method b) are all faint and/or blended. Results from method $b$ ) above, together with these corrections, are adopted for $\mathrm{C}, \mathrm{N}, \mathrm{O}$ abundances.

The results of our manual analysis differ from the outputs of the ASPCAP pipeline for oxygen and, to a lesser degree, nitrogen. Our methods a) and b) give rather similar results to each other within $\pm 0.05 \mathrm{dex}$, and with oxygen abundances somewhat higher than those derived with ASPCAP, that appear to be too low for bulge stars. The uncertainties on the oxygen abundances were already discussed by Jönsson et al. (2018) and Zasowski et al. (2019). Our oxygen abundances as compared with the DR17 ones are compatible within uncertainties, but with a trend to be higher.

In order to verify the reason for these differences, we carried out the fit to the N-rich star 2M17480576-2445000 analysed by Schiavon et al. (2017). With our method b) we have found that $[\mathrm{O} / \mathrm{Fe}]=0.4$ instead of $[\mathrm{O} / \mathrm{Fe}]=0.3$, and $[\mathrm{C} / \mathrm{Fe}]=-0.2$, instead of $[\mathrm{C} / \mathrm{Fe}]=0.0$, and on the other hand the N enhancement of $[\mathrm{N} / \mathrm{Fe}]=0.8$ is confirmed. Given the interplay between the CNO trio elements, it appears that the trend is to have somewhat lower C and higher O , and not much of a change in N abundances, in comparing our abundances with those from ASPCAP.

Note that none of the stars in our sample is N -enhanced, therefore they are good candidates to being similar to the first generation stars found in globular clustes.

## 3.2 alpha-elements $\mathrm{Mg}, \mathrm{Si}, \mathrm{Ca}$ and $\mathbf{T i}$

We analyse the abundances of the $\alpha$-elements $\mathrm{Mg}, \mathrm{Si}$, and Ca , and the iron-peak element Ti.

Magnesium, Silicon, Calcium and Titanium
Our fits with the original DR17 ASPCAP Mg abundances are in agreement with their results for $\mathrm{Mg}, \mathrm{Si}$ and Ca . Our calculations are in LTE with plane parallel models, as is adopted by the original ASPCAP method. The DR17 results for Mg and Ca correspond to calculations in non-LTE (Osorio et al. 2020), and even so the compatibility is good for these elements.

The SiI lines reported in Table 3 are all suitably reproduced with the ASPCAP Si abundance, with the exception of line SiI 15261.161 that is too shallow in the sample stars. Si abundance appears to be among the best determined ones by ASPCAP, together with Mg.

The four CaI lines listed in Smith et al. (2013) and that are use in ASPCAP, namely $16136.8,16150.8,16155.2,16157.4 \AA$ (see also Jönsson et al. 2018) are faint in the sample stars, and they are not fitted with the ASPCAP Ca abundance; instead they would need an extra 0.2 dex in Ca abundance to be fitted. In Table 3 we include another two lines of CaI that we were able to identify as suitable for the metallicity of our stars: CaI 16197.075, 16204.087 A. The CaI

Table 3. Line list. log gf from VALD3 linelist (Piskunov et al. 1995, Ryabchikova et al. 2015), Kurucz (1993) and NIST (Martin et al. 2002). The log gf values for CeII lines are from Cunha et al. (2017).

| Species | $\begin{gathered} \lambda \\ (\AA) \end{gathered}$ | $\begin{aligned} & \text { Xex } \\ & (\mathrm{eV}) \end{aligned}$ | $\begin{gathered} \log \mathrm{gf} \\ \text { (VALD3) } \end{gathered}$ | $\begin{gathered} \log g f \\ \text { (Kurucz) } \end{gathered}$ | $\begin{aligned} & \log \mathrm{gf} \\ & (\text { NIST } \end{aligned}$ | Notes |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| SiI | 15361.161 | 5.954 | -1.925 | -1.990 | -1.710 | Apogee gap |
|  | 15376.831 | 6.222 | -0.649 | -0.290 | - |  |
|  | 15833.602 | 6.222 | -0.168 | -0.660 | -0.078 |  |
|  | 15960.063 | 5.984 | 0.107 | 0.130 | 0.197 |  |
|  | 16060.009 | 5.954 | -0.566 | -0.440 | -0.429 |  |
|  | 16094.787 | 5.964 | -0.168 | -0.110 | -0.078 |  |
|  | 16215.670 | 5.964 | -0.665 | -0.990 | -0.575 |  |
|  | 16680.770 | 5.984 | -0.140 | -0.500 | -0.090 |  |
|  | 16828.159 | 5.984 | -1.102 | -1.390 | -1.012 |  |
| CaI | 16197.075 | 4.535 | 0.089 | 0.638 | - |  |
|  | 16204.087 | 4.535 | -0.627 | -0.111 | - |  |
| TiI | 15543.756 | 1.879 | -1.120 | -1.273 | -1.080 |  |
|  | 15602.842 | 2.267 | -1.643 | -1.544 | - |  |
|  | 15698.979 | 1.887 | -2.060 | -2.218 | -2.020 |  |
|  | 15715.573 | 1.873 | -1.250 | -1.359 | -1.200 |  |
|  | 16635.161 | 2.345 | -1.807 | -2.178 | - |  |
| CeII | 15277.610 | 0.609 | -1.94 | - | - | too faint |
|  | 15784.750 | 0.318 | -1.54 | - | - |  |
|  | 15829.830 | 0.320 | -1.80 | - | - | Apogee gap |
|  | 15958.400 | 0.470 | -1.71 | - | - |  |
|  | 15977.120 | 0.232 | -2.10 | - | - | weak line strongly blended |
|  | 16327.320 | 0.561 | -2.40 | - | - |  |
|  | 16376.480 | 0.122 | -1.79 | - | - |  |
|  | 16595.180 | 0.122 | -2.19 | - | - |  |
|  | 16722.510 | 0.470 | -1.65 | - | - |  |

Table 4. Solar abundances from (1) Grevesse et al. (1996, 1998) (adopted); (2) Steffen et al. (2015); (3) Scott et al. (2015a,b); (4) Grevesse et al. (2015); (5) Lodders et al. (2009); Arcturus abundances from: (6) Ramírez \& Allende Prieto (2011), (7) McWilliam et al. (2013), (8) Lecureur et al. (2007), (9) Barbuy et al. (2014), (10) Smith et al. (2013); (11) Cunha et al. (2017) $\mu$ Leo abundances from: (10) Smith et al. (2013); (12) Gratton \& Sneden (1990), 13: Barbuy et al. (2015), (14) Van der Swaelmen et al. (2016); (15) fits to the optical spectrum of $\mu$ Leo.

| El. | Z | $\begin{gathered} \log \epsilon(X)_{\odot} \\ \text { Sun } \end{gathered}$ |  |  | [X/Fe] $\quad \log \epsilon(X)$ <br> Arcturus adopted |  | $\begin{gathered} {[\mathrm{X} / \mathrm{Fe}] \underset{\substack{ \\ \mu \text { Leo } \\ \text { adopted }}}{\log \epsilon(X)}} \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Fe | 26 | 7.50 | 7.50 | 7.50 | -0.54 | 6.96 | +0.30 | 7.80 | 7.76 |
| C | 6 | 8.55[1] | - | 8.39[5] | -0.22[8] | 7.79 | -0.3[10] | 8.55 | 8.52 |
| N | 7 | 7.97[1] | - | 7.86[5] | +0.22[8] | 7.65 | +0.45[10] | 8.72 | 8.71 |
| O | 8 | 8.76[2] | - | 8.73[5] | 0.39 [9] | 8.62 | +0.0[10] | 9.06 | 9.05 |
| Na | 11 | 6.33[1] | 6.21[3] | 6.30[5] | 0.11[6] | 5.90 | +0.50[8] | 7.13 | - |
| Mg | 12 | 7.58[1] | 7.59[3] | 7.54[5] | 0.37[6] | 7.41 | -0.03[10] | 7.85 | 7.85 |
| Al | 13 | 6.47[1] | 6.43 [3] | 6.47[5] | 0.37[7] | 6.30 | +0.13[10] | 6.90 | 6.90 |
| Si | 14 | 7.55[1] | 7.51[3] | 7.52[5] | 0.33[6] | 7.34 | -0.10[10] | 7.75 | 7.76 |
| Ca | 20 | 6.36[1] | 6.32[3] | 6.33[5] | 0.11[6] | 5.93 | -0.04[10] | 6.62 | 6.62 |
| Sc | 21 | 3.17[1] | 3.16[3] | 3.10[5] | 0.23[6] | 2.86 | $+0.10[11]$ | 3.57 | - |
| Ti | 22 | 5.02[1] | 4.93[3] | 4.90[5] | 0.26[7] | 4.74 | +0.10[10] | 5.42 | 5.40 |
| V | 23 | 4.00[1] | 3.89[3] | 4.00[5] | 0.12[7] | 3.58 | +0.03[12] | 4.33 | 4.18 |
| Cr | 24 | 5.67[1] | 5.62[3] | 5.64[5] | -0.05[6] | 5.08 | -0.01[12] | 5.96 | 6.14 |
| Mn | 25 | 5.39[1] | 5.42[3] | 5.37[5] | -0.14[7] | 4.71 | $+0.00[12]$ | 5.69 | 5.79 |
| Co | 27 | 4.92[1] | 4.93[3] | 4.92[5] | +0.09[7] | 4.49 | +0.00[12] | 5.22 | 5.23 |
| Ni | 28 | 6.25[1] | 6.20[3] | 6.23[5] | 0.06[6] | 5.77 | +0.05[10] | 6.60 | 6.60 |
| Cu | 29 | 4.21[1] | 4.19[4] | 4.21[5] | -0.26[10] | 3.55 | -0.10[10] | 4.41 | 4.41 |
| Zn | 30 | 4.60[1] | 4.56[4] | 4.62[5] | +0.18[6] | 4.26 | -0.10[13] | 4.80 | - |
| Y | 39 | 2.24[1] | 2.21[4] | 2.21[5] | -0.30[9] | 1.40 | +0.04[14] | 2.58 | - |
| Zr | 40 | 2.60[1] | 2.59[4] | 2.58[5] | -0.28[7] | 1.78 | +0.10[12] | 3.00 | - |
| Ba | 56 | 2.13[1] | 2.25[4] | 2.17[5] | -0.30[9] | 1.29 | +0.10[14] | 2.53 | - |
| La | 57 | 1.22[1] | 1.11[4] | 1.14[5] | -0.30[9] | 0.38 | -0.37[14] | 1.15 | - |
| Ce | 58 | 1.55[1] | 1.58[4] | 1.61[5] | -0.45[11] | 0.99 | -0.37[14] | 1.15 | - |
| Eu | 63 | 0.51[1] | 0.52[4] | 0.52[5] | 0.23[7] | 0.20 | -0.14[14] | 0.67 | - |



Figure 4. Star $2 \mathrm{M} 17382504-2424163$ : Observed spectrum (black dotted) and synthetic spectrum computed with $[\mathrm{C}, \mathrm{N}, \mathrm{O} / \mathrm{Fe}]=-0.20,0.30,0.40(\mathrm{green})$.
16197.075 is well fitted in about half the stars, whereas in others it show blends, and finally the CaI $16204.087 \AA$ line is well fitted with the Ca abundance from ASPCAP. A FWHM $=0.65$ fits better the lines. In conclusion, we adopted the ASPCAP Ca abundances, relying on the results for the CaI 16204.087 A line.

Titanium: Among the 5 lines studied, only TiI 15543.756 Å line is well fit in essentially all stars. TiI $15698.979 \AA$ tends to give the same value, but it is located in a blend with several other lines, with a difficult continuum placement. TiI 15715.753 Å tends to give either the value from ASPCAP or requires a lower Ti abundance, whereas TiI 15602.842 , and $16635.161 \AA$ require higher values by about $0.3 \pm 0.2$ dex to be fitted. Because of the conflicting results from these different lines, and the fact that ASPCAP gives $[\mathrm{Ti} / \mathrm{Fe}]=0.0$ for most stars, which is not compatible with the Si and Ca enhancements, we preferred not to analyse the Ti abundances in the sample stars.

Note that the lines TiI 15602.842 , and $16635.161 \AA$, that are only fitted with higher Ti abundances, have somewhat higher excitation potential than the other 3 inspected lines, and that means that there may be an effect of effective temperature.

## 3.3 s-process element Ce

We used 6 CeII lines, among which CeII $16722.510 \AA$ line is well fit to almost all stars, except for a few for which most of the other lines are fit with a lower value than with the best line (case of 2M171736932806495), followed by CeII 15958.400 and $16595.180 \AA$ lines, that are fit with the adopted value for almost all stars.

CeII $15784.750 \AA$ is fit for about half the stars, for a few would require lower Ce abundances and about $1 / 3$ of them would require higher Ce abundances; $16327.320 \AA$ is faint and is fit for about $1 / 3$ of stars and $2 / 3$ would require higher Ce abundances; CeII 16376.480 $\AA$ would require higher values for about half the stars.

In 8 cases all six lines can be considered well-fitted, as is the case of star 2M18500307-1427291, shown in Figure 8.

For the fit of the Ce lines, we adopted $\mathrm{FWHM}=0.75$, which is suitable for the wavelength of the lines. The revised values are systematically higher than those resulting from ASPCAP (the fits to all stars are available under request).
DR17 used three Ce II windows covering the lines 15784, 16376, and $16595 \AA$. These are the stronger lines among the six that we used, and for all the sample stars it is clear that, from these 3 lines, a higher Ce abundance is needed.


Figure 5. Star 2M17382504-2424163: Selected OH lines. Observed spectrum (black dotted) and synthetic spectrum computed with $[\mathrm{O} / \mathrm{Fe}]=0.40$ (green).

We note that due to uncertainties in the Ce abundances in DR17 the APOGEE team has released internally to the collaboration a value added catalogue with revised abundances (Hayes et al. in preparation). This catalogue will be public to the community in a few months.

As for Nd we found that the lines are not suitable for analysis, from fits to them in the reference stars Arcturus and $\mu$ Leo, therefore we disregarded this element in the present analysis.

## 4 DISCUSSION

The $\alpha$-element abundances in bulge stars provide us with a constraint on the formation history of its stellar populations: the formation timescale. In other words, a mean $[\alpha / \mathrm{Fe}] \sim 0.5$ in halo and bulge metalpoor stars of $[\mathrm{Fe} / \mathrm{H}] \lesssim-1.0$ indicates a fast chemical enrichment at early times, dominated by supernovae type II (SNII) (e.g. Woosley \& Weaver 1995, hereafter WW95), whereas a lower [ $\alpha / \mathrm{Fe}$ ] implies a slower enrichment, allowing supernovae type Ia to contribute to the enrichment of iron.

Moreover, as recently shown by Miglio et al. (2021) for a sample
of Kepler stars with APOGEE spectra, stars with $[\alpha / \mathrm{Fe}]>0.2$ are all very old. The same probably applies to the present sample.

### 4.1 Oxygen and magnesium

Oxygen is produced by helium and neon burning in hydrostatic phases of the evolution of massive stars. Magnesium, together with Aluminum, are produced in hydrostatic carbon and neon burning (WW95). O and Mg are therefore the bona-fide alpha-elements produced by massive stars and ejected by supernova type II (SNII) event. They are enhanced relative to iron in old stellar populations, such as in bulge stars. In Figure 9 (upper panel) are plotted the oxygen abundances reported by the original APOGEE DR17 release, and the revised values obtained as explained in Section 3. This Figure is readapted from that in Barbuy et al. (2018a), taking into account only the literature higher-resolution data (with a few exceptions) and data showing little abundance spread. The literature data taken into account are from Friaça \& Barbuy (2017), that contains a revision of the abundances from Zoccali et al. (2006) and Lecureur et al. (2007), Cunha \& Smith (20060, Alves-Brito et al. (2010), Fulbright et al. (2007), only stars older than 11 Gyr from Bensby et al. (2013), Ryde et al. (2010) including a few of the same stars from Zoccali et al.


Figure 6. Stars 2M17382504-2424163 and 2M17511568-3249403: Selected CO lines. Observed spectra (black dotted) and synthetic spectra computed with $[\mathrm{C} / \mathrm{Fe}]=-0.2$ for both stars, and $[\mathrm{O} / \mathrm{Fe}]=0.40$ and 0.38 respectively (green).
(2006), Jönsson et al. (2017) including reanalysed 23 stars from Zoccali et al. (2006), Lecureur et al. (2007) and Friaça \& Barbuy (2017), Siqueira-Mello et al. (2016), and metal-poor stars from García-Pérez et al. (2013), Howes et al. (2016) and Lamb et al. (2017).

Figure 9 (lower panel) gives [ $\mathrm{Mg} / \mathrm{Fe}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$ for metal-poor stars from García-Pérez et al. (2013), Howes et al. (2016), Lamb et al. (2017), Casey \& Schlaufman (2015), Koch et al. (2016), Fulbright et al. (2007), as corrected by McWilliam (2016), Alves-Brito et al. (2010), Hill et al. (2011), Bensby et al. (2017) for stars older than 8 Gyr, Johnson et al. (2014), Ryde et al. (2010, 2016), Siqueira-Mello et al. (2016), Jönsson et al. (2017) and Rojas-Arriagada et al. (2017).

Our fits show agreemeent with the APOGEE ASPCAP Mg abundances, and they are compatible with the Mg abundances of other samples of bulge stars. The different model lines in Figure 9 correspond to different radii from the Galactic center.

### 4.2 Chemodynamical evolution model for O and Mg

We have compared the abundances derived from observations with the predictions of chemodynamical evolution models for the bulge (Friaça \& Barbuy 2017), described as a classical spheroid. It is
assumed a baryonic mass of $2 \times 10^{9} \mathrm{M}_{\odot}$, and a dark halo mass $M_{H}=$ $1.3 \times 10^{10} \mathrm{M}_{\odot}$. One central parameter of the model is the specific star formation rate $v_{S F}$ (i.e. the inverse of the star formation time scale).

In the nucleosynthesis prescriptions of our model, we adopt the metallicity dependent yields from core-collapse supernovae (SNe II) from WW95, with some modifications following suggestions of Timmes et al. (1995). For low metallicities ( $\mathrm{Z}<0.01 \mathrm{Z}_{\odot}$ ), we included the yields from high explosion-energy hypernovae (HNe) (Nomoto et al. 2013, and references therein). The type Ia supernovae yields are from Iwamoto et al. (1999) - their models W7 (progenitor star of initial metallicity $\mathrm{Z}=\mathrm{Z}_{\odot}$ ) and W 70 (zero initial metallicity). The yields for intermediate mass stars $\left(0.8-8 \mathrm{M}_{\odot}\right)$ ) with initial $\mathrm{Z}=0.001$, $0.004,0.008,0.02$, and 0.4 come from van den Hoek \& Groenewegen (1997) (variable $\eta_{A G B}$ case).

As we can see from Figure 9, the abundances derived here both for the oxygen (upper panel) and for the magnesium (lower panel) are well reproduced by the chemodynamical model with $v_{S F}=1 \mathrm{Gyr}^{-1}$ (star formation time scale of 1 Gyr ). Once more this suggests these objects to be very old.


Figure 7. Star 2M17511568-3249403: regions containing CN lines. Observed spectra (black dotted) and synthetic spectra (green) computed with [C/Fe] = -0.2, $[\mathrm{N} / \mathrm{Fe}]=0.0,[\mathrm{O} / \mathrm{Fe}]=0.38$; green dotted lines correspond to calculations with $[\mathrm{N} / \mathrm{Fe}]=0.2,0.4$.

### 4.3 Silicon and calcium

Si and Ca are mainly produced during the explosive nucleosynthesis of SNII events (WW95; McWilliam 2016), with smaller contributions from supernovae type Ia (SNIa).

The $\alpha$-elements Si and Ca are plotted in Figure 10 for the 58 sample stars together with literature data from García-Pérez et al. (2013), Howes et al. (2016), Lamb et al. (2017), Casey \& Schlaufman (2015), Koch et al. (2016), Alves-Brito et al. (2010), Bensby et al. (2017) for stars older than 8 Gyr, Ryde et al. (2016) and SiqueiraMello et al. (2016).

Figure 10 shows that a typical star formation time scale of 1 Gyr (the chemodynamical model with $v_{S F}=1 \mathrm{Gyr}^{-1}$ ) also explains the Si and Ca abundances found in the bulge.

From this figure we can conclude that there are no differences in the Si and Ca abundances of the present confirmed in-situ samples of bulge stars, and previous samples in bulge regions, for which no precise distances were available.

### 4.4 Heavy element Ce

In Figure 11 are shown the revised Ce abundances in contrast with the lower original DR17 APOGEE abundances, together with results for M62 from Yong et al. (2014) and for field bulge stars from van der Swaelmen et al. (2016) and Lucey et al. (2022). As can be seen from the figure, we have found that the sample stars are enhanced in Cerium, by about a mean value of $[\mathrm{Ce} / \mathrm{Fe}] \sim 0.4$. We find Ce enhancements relative to the DR17 and as well to results from the COMBS survey by Lucey et al. (2022). Cleary further investigation on Ce abundances in metal-poor bulge stars is needed.
$\mathrm{Ce}(\mathrm{Z}=58, \mathrm{~A}=140)$ is essentially an element mostly formed by the s-process, with a fraction of 0.186 as r-element and 0.814 as s-element (Simmerer et al. 2004). Ce appears to be overproduced in massive spinstars (Frischknecht et al. 2016). On the other hand, since we do not have the ages of these stars, we cannot exclude that the Ce enhancement could be due to a mass transfer from a companion Asymptotic Giant Branch (AGB) star (e.g Bisterzo et al. 2011, Cristallo et al. 2015).

Table 5. Revised CNO abundances derived from non-calibrated DR17 stellar parameters compared with the DR17 CNO abundances in the last column.



Figure 8. Star 2M18500307-1427291: fit to the 6 cerium lines. Observed spectrum: black; śynthetic spectra are: blue with original ASPCAP [Ce/Fe]=0.08 Ce abundance, green with final Ce abundance.

## 5 CONCLUSIONS

We have selected 58 stars from the bulge sample of Queiroz et al. (2021) based on APOGEE, with characteristics to belong to the spheroidal, pressure-supported bulge. For this sample we have analysed lines of C, N, O, alpha-elements $\mathrm{Mg}, \mathrm{Si}, \mathrm{Ca}$ and neutron-capture element Ce . Our fits to the $\mathrm{Mg}, \mathrm{Si}$ and Ca abundances from the original APOGEE results using the ASPCAP software appeared to be in good agreement. We recomputed abundances for $\mathrm{C}, \mathrm{N}, \mathrm{O}$, and Ce , assuming the spectroscopic non-calibrated stellar parameters from APOGEE DR17. We report differences in abundances of these elements.

We compare the abundances of these elements to literature data for bulge stars, and chemodynamical models by Friaça \& Barbuy
(2017) - see also Barbuy et al. (2018a). These comparisons show compatibility of the abundances of the sample stars with literature and models for $\mathrm{Mg}, \mathrm{Si}$, and Ca in which a pressure supported component (spheroidal bulge) formed on a very short timescale (below 1 Gyr ). Similar results were suggested by other chemical evolution models (see Matteucci 2021 for a review), and for stars with similar alpha element enhancements with asteroseismic ages (Miglio et al. 2021).

Nitrogen abundances show no exceptional enhancement for any of the sample stars, therefore there is no evidence for these stars to be a result of multiple stellar populations in dissolved globular clusters. The Ce abundance is enhanced in all stars, which would point out to a s-process origin of this element already in the very early phases of chemical enrichment. This could have been achieved with spinstars


Figure 9. $[\mathrm{O} / \mathrm{Fe}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$ (upper panel) and $[\mathrm{Mg} / \mathrm{Fe}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$ (lower panel), for literature bulge field stars and the APOGEE abundances (original and revised in the case of oxygen) for the 58 sample stars. Symbols: grey 4-pointed stars: Alves-Brito et al. (2010); red filled circles: Bensby et al. (2013); red filled circles: Bensby et al. (2017); grey open pentagons: Casey \& Schlaufman (2015); strong-grey filled triangles: Fulbright et al. (2007); magenta open pentagons: García-Pérez et al. (2013); red filled squares: Cunha \& Smith (2006); indianred filled circles: Hill et al. (2011); grey open pentagons: Howes et al. (2016); grey stars: Johnson et al. (2014); grey 5-pointed stars: Jönsson et al. (2017); grey open pentagons: Koch et al. (2016); green open pentagons: Lamb et al. (2017); red crosses: Ryde et al. (2010); green filled circles: Ryde et al. (2016); turquoise 5-pointed stars: Rojas- Arriagada et al. (2017); grey open triangles: Siqueira-Mello et al. (2016); blue open circles: APOGEE original; cyan filled circles: final abundances for the 58 APOGEE sample stars. The oxygen abundances are normalized in terms of adopted solar abundances as explained in Friaça \& Barbuy (2017). Chemodynamical evolution models from Friaça \& Barbuy (2017) with formation timescale of 1 Gyr , for several radii, are overplotted: $r<0.5 \mathrm{kpc}$ (solid lines), $0.5<1 \mathrm{kpc}$ (dotted lines), $1<r<2 \mathrm{kpc}$ (short-dashed lines), $2<r<3 \mathrm{kpc}$ (long-dashed lines).
(e.g. Chiappini et al. 2011, Frischknecht et al. 2016), or alternatively due to mass transfer from a companion AGB star (e.g. Cristallo et al. 2015). This same conclusion was reached by Barbuy et al. (2009, 2014, 2021b) regarding the globular cluster NGC 6522, but here, since all the present sample stars are enhanced in Ce , all of them
would have to be binaries with an AGB companion. Therefore the spinstars seem to be a more plausible explanation.


Figure 10. $[\mathrm{Si}, \mathrm{Ca} / \mathrm{Fe}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$ for literature bulge field stars and the APOGEE abundances for the 58 sample stars. Symbols: grey 4-pointed stars: Alves-Brito et al. (2010); red filled circles: Bensby et al. (2017); grey open pentagons: Casey \& Schlaufman (2015); magenta open pentagons: García-Pérez et al. (2013); grey open pentagons: Howes et al. (2016); grey open pentagons: Koch et al. (2016); green open pentagons: Lamb et al. (2017); green filled circles: Ryde et al. (2016); grey open triangles: Siqueira-Mello et al. (2016); blue open circles: APOGEE original; cyan filled circles: final abundances for the 58 APOGEE sample stars. The lines are the predictions of the chemodynamical models of Friaça \& Barbuy (2017) with a formation timescale of 1 Gyr for several radii.

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Table 6. Abundances from original APOGEE-ASPCAP derivations, and revised values of Ce , using the lines reported in Table 3 for the 58 sample stars. For Ce abundances the two columns correspond to DR17 ASPCAP abundances, and revised values.

| Star | [ $\mathrm{Ce} / \mathrm{Fe}$ ] |  |
| :---: | :---: | :---: |
| ID | DR17 | revised |
| 2M17153858-2759467 | -0.16 | 0.25 |
| 2M17173693-2806495 | -0.10 | 0.20 |
| 2M17250290-2800385 | -- | 0.20 |
| 2M17265563-2813558 | -0.20 | 0.10 |
| 2M17281191-2831393 | -0.14 | 0.20 |
| 2M17295481-2051262 | -0.02 | -0.02 |
| 2M17303581-2354453 | -- | 0.40 |
| 2M17324257-2301417 | -- | - |
| 2M17330695-2302130 | -- | 0.50 |
| 2M17344841-4540171 | -- | 0.50 |
| 2M17351981-1948329 | -- | 0.50 |
| 2M17354093-1716200 | -- | 0.50 |
| 2M17390801-2331379 | -- | 0.50 |
| 2M17392719-2310311 | -- | 0.50 |
| 2M17473299-2258254 | -0.27 | 0.30 |
| 2M17482995-2305299 | -0.4 | 0.20 |
| 2M17483633-2242483 | -- | 0.50 |
| 2M17503263-3654102 | -- | 0.50 |
| 2M17552744-3228019 | -0.15 | 0.35 |
| 2M18020063-1814495 | -0.08 | 0.30 |
| 2M18050452-3249149 | -0.11 | 0.45 |
| 2M18050663-3005419 | -- | 0.40 |
| 2M18065321-2524392 | -- | 0.45 |
| 2M18104496-2719514 | -0.17 | 0.25 |
| 2M18125718-2732215 | -- | 0.30 |
| 2M18200365-3224168 | 0.07 | 0.50 |
| 2M18500307-1427291 | 0.08 | 0.45 |
| 2M17173248-2518529 | 0.03 | 0.45 |
| 2M17285088-2855427 | -- | 0.50 |
| 2M17291778-2602468 | -- | 0.45 |
| 2M17301495-2337002 | -- | 0.45 |
| 2M17310874-2956542 | -0.17 | 0.30 |
| 2M17382504-2424163 | -- | 0.10 |
| 2M17453659-2309130 | -0.40 | -0.10 |
| 2M17511568-3249403 | -0.11 | 0.40 |
| 2M17532599-2053304 | -- | 0.40 |
| 2M17552681-3334272 | 0.03 | 0.35 |
| 2M18005152-2916576 | -0.12 | 0.30 |
| 2M18010424-3126158 | -- | 0.43 |
| 2M18042687-2928348 | -0.07 | 0.20 |
| 2M18044663-3132174 | -- | 0.33 |
| 2M18052388-2953056 | -0.29 | 0.20 |
| 2M18080306-3125381 | 0.15 | 0.25 |
| 2M18142265-0904155 | -0.15 | 0.30 |
| 2M18195859-1912513 | -0.17 | 0.35 |
| 2M17190320-2857321 | -0.20 | 0.32 |
| 2M17224443-2343053 | 0.13 | 0.55 |
| 2M17292082-2126433 | -0.12 | 0.42 |
| 2M17293482-2741164 | -0.27 | 0.30 |
| 2M17323787-2023013 | -- | 0.42 |
| 2M17330730-2407378 | -0.18 | 0.30 |
| 2M17341796-3905103 | -0.03 | 0.20 |
| 2M17342067-3902066 | -0.18 | 0.20 |
| 2M17503065-2313234 | - | 0.20 |
| 2M18023156-2834451 | - | 0.50 |
| 2M18143710-2650147 | -0.18 | 0.20 |
| 2M18150516-2708486 | - | 0.25 |
| 2M18344461-2415140 | -0.28 | 0.40 |



Figure 11. $[\mathrm{Ce} / \mathrm{Fe}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$ for literature bulge field stars and the APOGEE abundances for the 58 sample stars. Symbols: red filled triangles: van der Swaelmen et al. (2016); green filled triangles: Lucey et al. (2022); magenta pentagons: M62 from Yong et al. (2014); open blue circles: APOGEE DR17 [Ce/Fe] values and cyan filled circles: revised abundances for the 58 APOGEE sample stars. Error bars are indicated in the right upper corner.

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## DATA AVAILABILITY

The observed data are from the APOGEE survey. The calculations and fits to the lines can be requested to the authors. The code for spectrum synthesis is available to be retrieved.

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[^2]:    1 Value added catalogues are available in both releases
    2 To estimate this probability, we used the Monte Carlo sample of each star ( 50 orbits) and calculated the fraction of orbits classified as bar-shaped

[^3]:    ${ }^{3}$ The code is available at http://trevisanj.github.io/PFANT.

